

**SUPPLEMENTARY INFORMATION**  
**Photoemission study of twisted monolayers and bilayers of WSe<sub>2</sub> on graphite substrates**

Bharti Parashar,<sup>1</sup> Lars Rathmann,<sup>2</sup> Hyun-Jung Kim,<sup>3</sup> Iulia Cojocariu,<sup>1</sup> Aaron Bostwick,<sup>4</sup> Chris Jozwiak,<sup>4</sup> Eli Rotenberg,<sup>4</sup> José Avila,<sup>5</sup> Pavel Dudin,<sup>5</sup> Vitaliy Feyer,<sup>1</sup> Bernd Beschoten,<sup>2,6</sup> Christoph Stampfer,<sup>2</sup> Gustav Bihlmayer,<sup>3</sup> Claus M. Schneider,<sup>1,7,8</sup> and Lukasz Plucinski<sup>1,\*</sup>

<sup>1</sup>*Peter Grünberg Institut (PGI-6), Forschungszentrum Jülich GmbH, 52428 Jülich, Germany*

<sup>2</sup>*2nd Institute of Physics and JARA-FIT and, RWTH Aachen University, 52074 Aachen, Germany*

<sup>3</sup>*Peter Grünberg Institut (PGI-1), Forschungszentrum Jülich GmbH, 52428 Jülich, Germany*

<sup>4</sup>*Advanced Light Source, Lawrence Berkeley National Laboratory,  
One Cyclotron Road, Berkeley, CA 94720, USA*

<sup>5</sup>*Synchrotron-SOLEIL, Université Paris-Saclay, Saint-Aubin, BP48, F91192 Gif sur Yvette, France*

<sup>6</sup>*JARA-FIT Institute for Quantum Information, Forschungszentrum  
Jülich GmbH and RWTH Aachen University, 52074 Aachen, Germany*

<sup>7</sup>*Fakultät für Physik, Universität Duisburg-Essen, 47048 Duisburg, Germany*

<sup>8</sup>*Physics Department, University of California, Davis, CA 95616, USA*

(Dated: December 18, 2022)

---

\* l.plucinski@fz-juelich.de

## SI. IDENTIFICATION OF SL AND BL MICRO-REGIONS ON SAMPLE *B*

The optical micrograph of sample *B* is presented in Fig. S1(a). We acquired a spatial map of ARPES intensity using capillary micro-focusing at  $h\nu = 147$  eV at the MAESTRO beamline, as shown in Fig. S1(b). The numbered white boxes in Fig. S1(b) indicate the direction of the line-scan spectra and the positions at which we collected the data set.

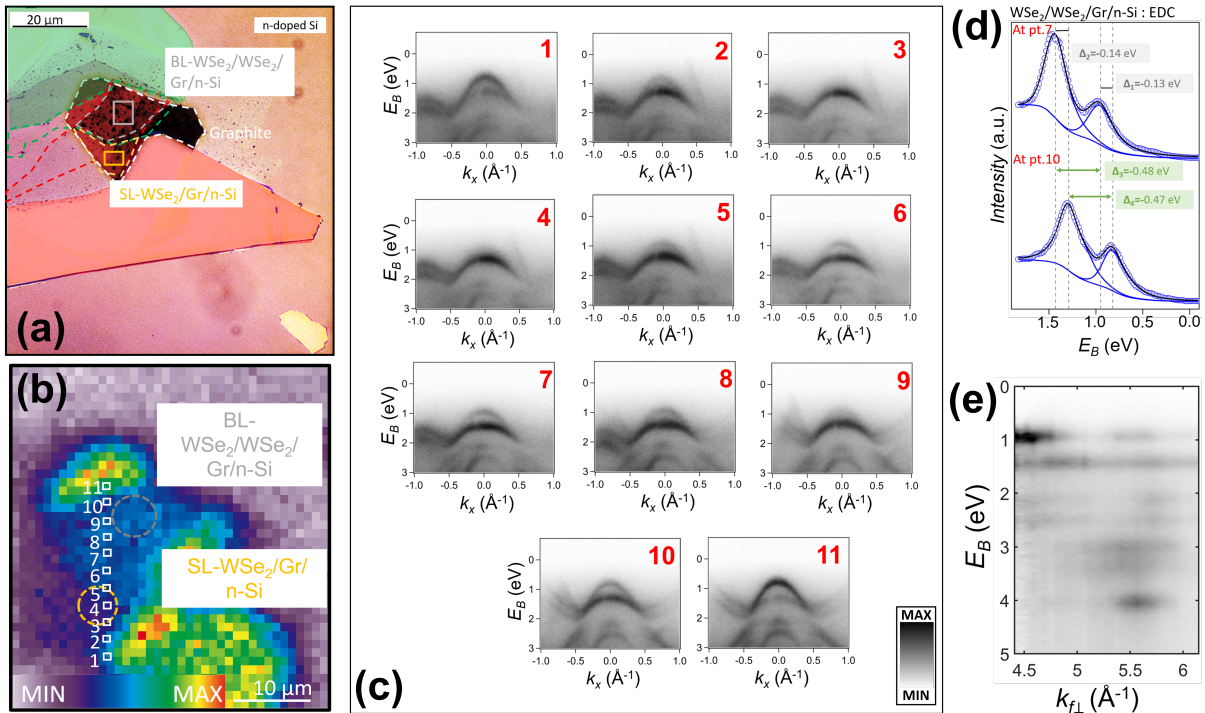


Figure S1. Characterization of sample *B*. (a) Optical micrograph of the twisted SL and BL WSe<sub>2</sub>/graphite/*n*-Si. (b) Spatial photoemission intensity map of an integrated intensity of the valence band at the momenta related to WSe<sub>2</sub>. (c) Collection of normal emission ARPES spectra at white boxes indicated in (b). (d) EDC profiles at boxes 7 and 10. (e) Set of normal emission spectra from twisted BL WSe<sub>2</sub>/graphite for  $h\nu$  between 80 and 150 eV with 2 eV step, with perpendicular momentum calculated as  $k_{f\perp} = \sqrt{(E_{kin} + V_0)2m/\hbar^2}$ , and  $V_0$  set to 10 eV.

Figure S1(c) depicts a set of ARPES maps at 11 boxes indicated in panel Figure S1(b).  $E(k)$  maps measured at normal emission at boxes 1, 6, 7, 8, 10, 11 show double bands and at boxes 3, 4 show single band near the VBM, which are characteristics signatures of the bilayer and monolayer regions, respectively. This indicates that also in sample *B* we were able to exclusively probe either SL or BL areas by changing the location of the beam position on the sample surface. For the BL regions, we measured binding energy shifts of the order of  $\sim 130$  meV between points 7 and 10. The peak positions are confirmed by the Voigt fit of normal emission EDCs, Fig. S1(d). This indicates inhomogeneities and local doping on the flake surface that vary on the  $\mu\text{m}$  scale. Figure S1(e) presents the photon energy dependence  $E_B(k_{\perp})$  map on sample *B* showing virtually the same features as sample *A* (Fig. 1(f)), even though both samples have different twist angles between the WSe<sub>2</sub> monolayers.

## SII. NANOARPES MEASUREMENTS AT ANTARES

Sample *A* had been characterized at the ANTARES beamline at SOLEIL synchrotron. As seen in Fig. S2(a), utilizing nanoARPES in the imaging mode, different regions of the WSe<sub>2</sub> flake can be distinguished. The electronic band structure could be acquired using the nanoARPES beamline in spectroscopic mode, as shown in Fig. S2(b,c), which depict the band dispersion of monolayer and bilayer regions. At the high symmetry  $\Gamma_0$  point, the monolayer region exhibits a single band at the local valence band maximum. In the bilayer region, the band is split into two branches. The observed results at ANTARES beamline are in agreement with those presented in the main paper, obtained at the MAESTRO beamline.

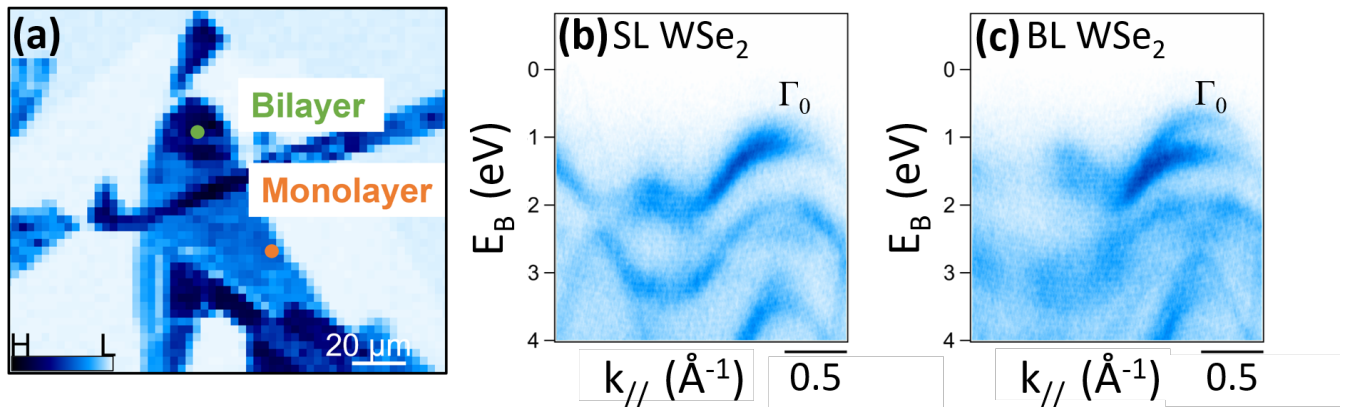


Figure S2. Band structure of different regions of WSe<sub>2</sub> flake by nanoARPES at ANTARES/SOLEIL. Panel (a) shows the nanoARPES image obtained by integrating intensity from 0.5 eV to 0.9 eV. Panels (b,c) show the band dispersion taken at  $h\nu = 100$  eV along the  $\Gamma$ - $K$  direction of the Brillouin zone from the two regions labeled in panel (a) by green and orange dots.

### SIII. SAMPLE C: TWISTED SL WSe<sub>2</sub>/GRAPHITE

Sample *C* was fabricated by mechanical exfoliation and dry transfer techniques. The WSe<sub>2</sub> monolayer was exfoliated from bulk crystal and transferred such that it partially covered two different substrates, graphite and *n*-Si. The overlapping area between WSe<sub>2</sub> layers and graphite is larger than  $20 \times 20 \mu\text{m}^2$  as shown in Fig. S3.

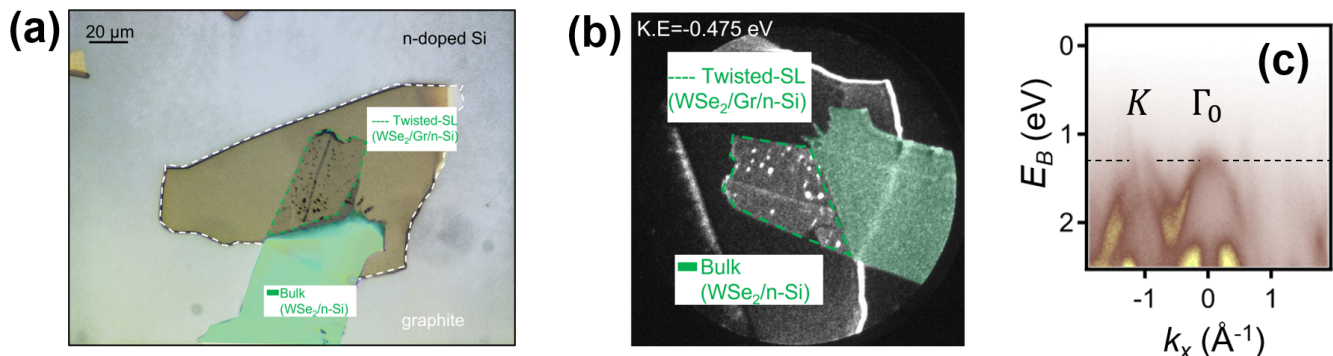


Figure S3. Overview of sample *C*. (a) Optical micrograph of twisted SL WSe<sub>2</sub> placed on graphite with the entire sample structure lying on the top of *n*-Si. (b) PEEM image of the sample taken with a Hg lamp at NanoESCA beamline at Elettra, where twisted SL and bulk WSe<sub>2</sub> flakes are represented by the dotted green lines and shaded green color regions, respectively. (c)  $E(k_x)$  map through  $\Gamma_0$  indicating single band at normal emission.

### SIV. ENERGY DISTRIBUTION CURVES AT $\Gamma_0$ AND $K$ IN TWISTED BL WSe<sub>2</sub>

To determine the band position from the ARPES data, we plotted the EDCs of the valence band dispersion at high symmetry  $\Gamma_0$  and  $K$  points of twisted WSe<sub>2</sub> of sample *B*. The particular EDCs in Fig. S4(a) and (b) are related to Figs. 6 (b) and (f) of the main manuscript, respectively. At  $\Gamma_0$  and  $K$  points in BL WSe<sub>2</sub>/graphite of sample *B* the local VB lie respectively at  $E_B \simeq 0.95$  eV and  $\simeq 0.72$  eV, indicating that the global VBM in twisted WSe<sub>2</sub> BL resides at  $K$  point. A splitting between the topmost  $K$  bands is approx. 480 meV.

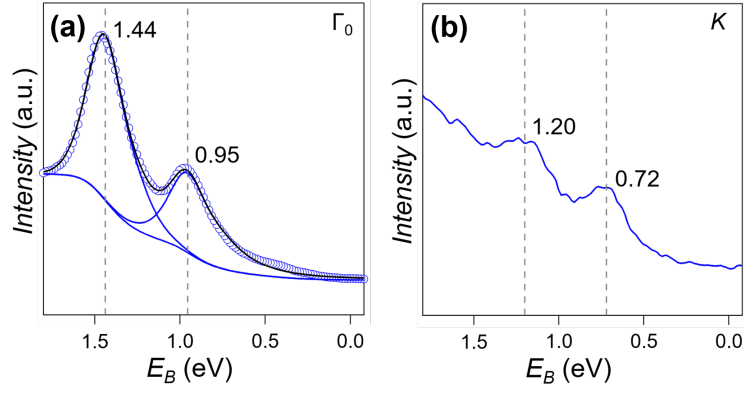


Figure S4. EDCs of twisted BL WSe<sub>2</sub> of sample *B* at  $\Gamma_0$  and  $K$ .

### SV. DETERMINATION OF THE TWIST ANGLE IN BL-WSE<sub>2</sub>

ARPES studies are often limited to mapping a limited portion of the surface Brillouin zone (SBZ), and probing dispersions simultaneously in both the first and the second SBZ is not performed routinely. Since the parallel momentum  $k_{\parallel}$  depends on the emission angle  $\theta$  and kinetic energy  $E_{kin}$  as  $k_{\parallel} = \sqrt{\frac{2m}{\hbar^2} E_{kin}} \sin \theta$ , there are often technical limitations in reaching large enough  $\theta$  at low  $E_{kin}$ , or a sufficient resolution at larger  $h\nu$ . The size of the SBZ depends on the lattice constant  $a$  as  $2\pi/a$ ; therefore, for many materials, one needs to reach  $k_{\parallel} \simeq 2 \text{ \AA}^{-1}$  to map the second SBZ, and this is, for example, out of reach using a helium discharge lamp with the He I line at  $h\nu \simeq 21.2 \text{ eV}$ . Also momentum microscopes exhibit a limit of the maximum parallel momentum at higher  $E_{kin}$  [1], which depends on the high voltage applied between the sample and the spectrometer, which often must be limited to control the field emission. The flexibility of MAESTRO beamline allows one to use  $h\nu = 147 \text{ eV}$  for 3D  $I(E_{bin}, k_x, k_y)$  mapping, allowing us to reach the second SBZ of WSe<sub>2</sub> at moderate emission angles of  $\theta \simeq 20^\circ$ .

Figure S5 shows constant energy surfaces of twisted BL WSe<sub>2</sub>/graphite of sample *A* and *B* across first and second SBZs at  $E_B \simeq 0.9 \text{ eV}$ , where we were able to visualize the intensity pockets at  $\Gamma_0$  and  $\Gamma_1$  contributing from both WSe<sub>2</sub> monolayers that constitute the BL WSe<sub>2</sub> structures. The characteristic features are labeled by the white ( $\Gamma_0$ ), red (top flake,  $\Gamma_{1top}$ ), and green (bottom flake,  $\Gamma_{1bottom}$ ) dotted color circles, as illustrated in Fig. S5.

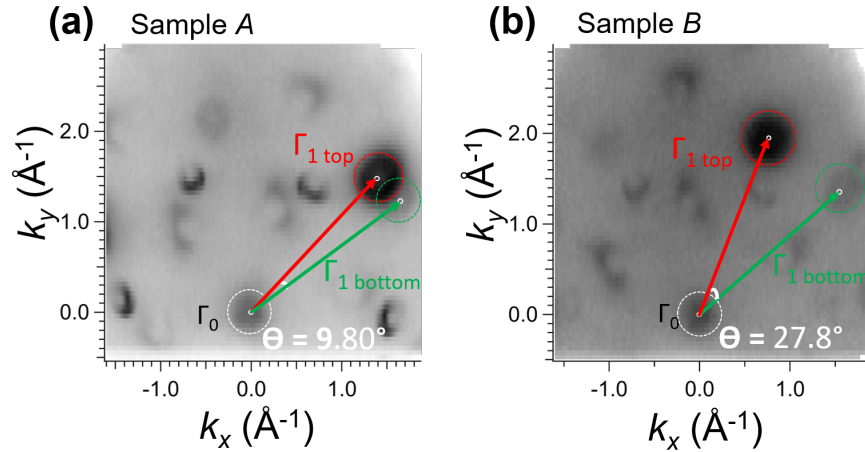


Figure S5. Panels (a) and (b) present the  $(k_x, k_y)$ -dependent constant energy maps taken at  $E_B \simeq 0.9 \text{ eV}$  of samples *A* and *B*, respectively. The red and green arrows connect  $\Gamma_0$  and  $\Gamma_{1top}$ ,  $\Gamma_0$  and  $\Gamma_{1bottom}$  high symmetry points of WSe<sub>2</sub> bilayer structures indicating the twist angle between BL WSe<sub>2</sub>.

Through identification of characteristic features from the top and bottom monolayers, we were able to determine the twist angles between the BL WSe<sub>2</sub> of sample *A* and *B*. For doing so, we utilize the band contours around the high symmetry  $\Gamma$  points on CEMs as illustrated in Figs. S5, S6, and S7. We mapped out the 3D  $I(E_B, k_x, k_y)$

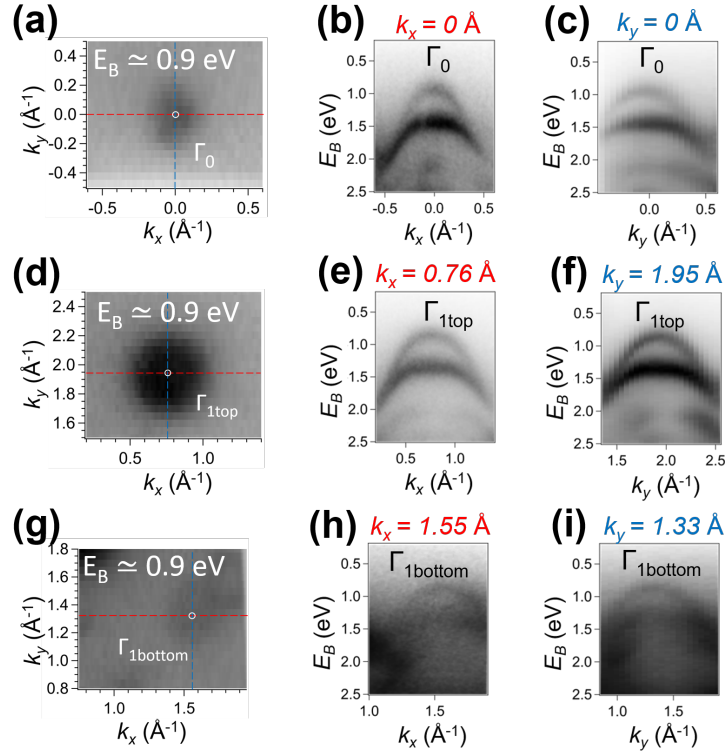


Figure S6. Sample *B*: (a,d,g) Constant energy contours at  $E_B \simeq 0.9$  eV taken around  $\Gamma_0$ ,  $\Gamma_{1top}$  and  $\Gamma_{1bottom}$  respectively. (b,c,e,f,h,i) Extracted  $E(k_x)$ ,  $E(k_y)$  dispersions at different  $k_x$  and  $k_y$  cuts performed on all gammas, as indicated by dashed red and blue lines in (a,d,g). A white circle marking in the (a,d,g) indicates the  $\Gamma_0$ ,  $\Gamma_{1top}$ , and  $\Gamma_{1bottom}$  positions retrieved from the 3D dataset.

dispersion of each bilayer structures and their respective monolayers. The vectors connecting  $\Gamma_0$  with  $\Gamma_{1top}$  and with  $\Gamma_{1bottom}$  are determined by identifying the centers of  $\Gamma$  pockets, as illustrated in Figs. S6, and S7. The selection of the  $k$  coordinates of the high symmetry  $\Gamma$  points is based on a comprehensive examination of the complete  $(E_B, k_x, k_y)$ -dependent intensity of BL WSe<sub>2</sub> structures. We analyzed  $E(k_x)$  and  $E(k_y)$  dispersions across all the  $\Gamma$  pockets originating from top and bottom monolayers features in BL WSe<sub>2</sub> datasets of sample *A* and *B* as shown in Figs. S6 and S7. We used all the  $\Gamma$  pockets visualized on the CEM from the WSe<sub>2</sub> bilayers datasets. The intensity of the  $\Gamma_{1bottom}$  is relatively low in sample *A*, making it difficult to select the  $k$  coordinates. For this, we approximated the coordinates from the monolayer dataset. The angle between the red and green vectors corresponds to the twist angles, which we determine as  $9.80^\circ \pm 1^\circ$  and  $27.80^\circ \pm 1^\circ$  of sample *A* and *B* as shown in Fig. S5. We note that due to non-linearities in the angular scale of the hemispherical analyzer these values very likely differ from the real twist angles, as discussed in the main text.

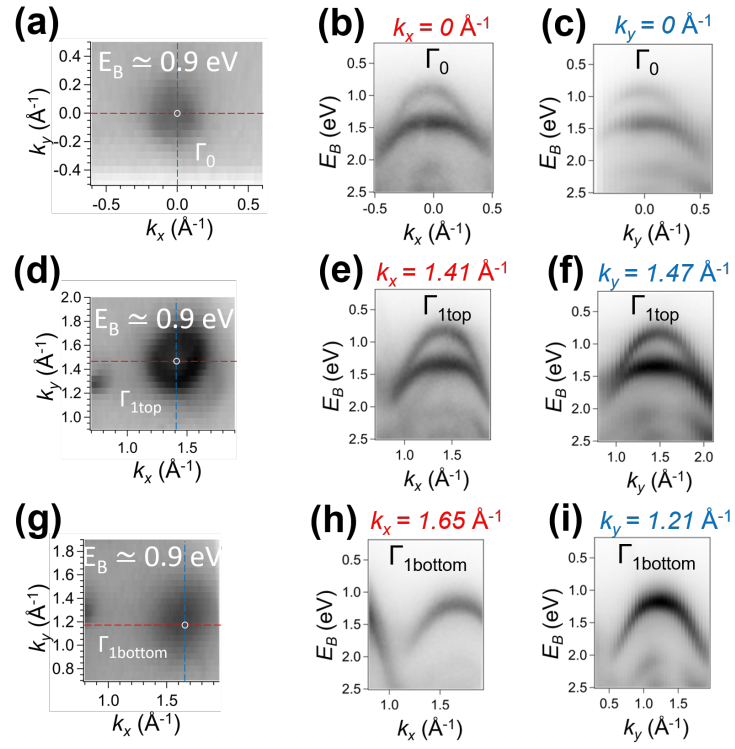


Figure S7. Same as Fig. S6 but for sample A.

- 
- [1] C. Tusche, A. Krasnyuk, and J. Kirschner, Spin resolved bandstructure imaging with a high resolution momentum microscope, *Ultramicroscopy* **159**, 520 (2015).