

# The Interaction of Sub-Monolayer Ta Adatoms and Clusters with Oxygen at the Pt(111) Interface

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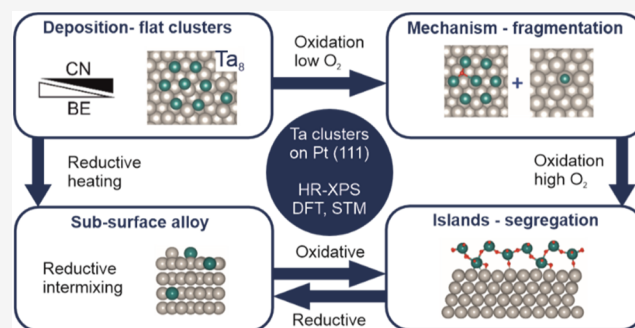


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**ABSTRACT:** The interaction of submonolayer quantities of size-selected and soft-landed Ta<sub>n</sub> (*n* = 4, 5, 6, 8, 13) clusters with Pt(111) is investigated employing high-resolution X-ray photoelectron spectroscopy (HR-XPS), scanning tunneling microscopy (STM), and density functional theory (DFT) simulations. The deposited clusters are monodispersed and stable under ultrahigh vacuum (UHV) conditions at 40 K. They display a size-specific trend in photoemission spectra, which is reasoned in terms of the distinct in plane coordination of Ta atoms in the clusters. Both the Ta coordination number and distance from the Pt surface influence its Bader charge and, accordingly, the oxidation state of the atoms in the Ta cluster. They already fragment in the presence of low amounts of oxygen and form a common oxidation product observed for all cluster sizes. Based on our observations, we propose an oxidation mechanism in the example of Ta<sub>8</sub> clusters, which is closely comparable to the one discussed in gas-phase studies on the oxidation of cationic Ta clusters of similar size. Concomitant to oxidation-induced fragmentation, the agglomeration into Ta-oxide islands with Ta in an oxidation state of +5 is observed. However, the strong interaction with the Pt surface leads to Ta 4f orbital photoemission features that differ from those commonly observed for Ta<sub>2</sub>O<sub>5</sub>. Computational insights concerning the structure of the Ta-oxide islands indicate flat agglomerates that agree with STM observations. They suggest distinct Ta 4f photoemission contributions from interfacial and surface-related Ta configurations. The respective HR-XPS spectra display specific core-level shifts as a function of bonding configuration and vicinity to the Pt surface. By annealing at 900 K in UHV, we observe oxygen loss and concomitant intermixing of Ta atoms with the Pt subsurface lattice to which results in the formation of a Ta–Pt alloy. These species, Ta-oxide islands, and Ta–Pt alloy, can reversibly interconvert by oxidative surface segregation and reductive intermixing.



## INTRODUCTION

Ta and Ta-oxide film-coated Pt components are widely employed in electronic devices, e.g., multilayer ceramic capacitors.<sup>1,2</sup> Coating the noble metal Pt with a relatively inert Ta-oxide of high dielectricity offers several advantages. The main benefit is increased chemical durability, such as of Ta<sub>2</sub>O<sub>5</sub>-coated Pt electrodes in electrochemical processes.<sup>3–6</sup> Besides an enhanced performance, new functionalities can emerge from particular arrangements that can be exploited to engineer novel applications, such as Schottky diodes for hydrogen sensing<sup>7</sup> and memristors for resistive random-access memory applications.<sup>8–12</sup> In these, (i) the oxidation state of Ta, (ii) the oxygen stoichiometry, and (iii) the migration of constituents are decisive for functionality.<sup>13–15</sup> Consequently, studying the electronic properties of interfacial structures constituted of Pt, Ta and Ta-oxides can yield fundamental

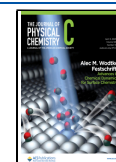
insights into their structure and evolution in different environments, which are of great interest both from a fundamental and applied point of view. The chemical interaction between Ta and Pt in Ta–Pt and Ta–O–Pt interfaces and its evolution during oxidation and film growth remains poorly understood. In particular, an unambiguous attribution of Ta oxidation states remains challenging mainly due to the numerous oxides and suboxides Ta can form and due to the electronic perturbation Ta experiences at the

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interface to Pt. Thus, attributing a particular oxidation state of Ta to a core-level shift (CLS) observed in X-ray photoelectron spectroscopy (XPS) is often inconclusive, and considerable discrepancies can be found in the literature.<sup>16–19</sup> These mainly emerge from the predominant technique used to investigate Ta-oxide thin films, sputter-etching combined with XPS.<sup>6,16,18,20–22</sup> While the method is straightforward and offers valuable insights (both quantitatively and qualitatively) into the chemical composition of a Ta-oxide film, it suffers from limitations due to its invasive nature. The bombardment with high-energy ions (in the kV range) inevitably induces changes in the original depth distribution, and computational reconstruction methods are often required to access information on the unperturbed structure.<sup>23–26</sup> Hence, pristine structural and chemical information is hardly discernible or not accessible, especially considering the interfacial contributions.

In contrast to sputter-etching, we present a bottom-up study of single Ta atoms and clusters deposited on a Pt(111) single crystal, shedding light on their structural, electronic, and chemical characteristics while emphasizing the interfacial constitution. We investigate the behavior of metallic Ta clusters at the Pt(111) surface and their reaction with oxygen at low (40 K) and elevated (900 K) temperatures using high-resolution X-ray photoelectron spectroscopy (HR-XPS) complemented by density functional theory (DFT) simulations and scanning tunneling microscopy (STM). Hereby, the oxidation of Ta atoms and clusters is monitored upon increasing exposure to O<sub>2</sub>, elucidating the oxidation products.

In previous work, we recently investigated the interaction of evaporated Ta atoms with the Pt(111) surface at 40 K.<sup>27</sup> Different surface adsorption sites for Ta adatoms were identified, with preferred adsorption to step edges. Summarizing this article, we found that the transient mobility of the evaporated atoms, with a residual kinetic energy of 0.28 eV/atom, promotes the formation of small Ta agglomerates. At elevated temperatures (900 K) in ultrahigh vacuum (UHV), Ta diffusion into the Pt surface is enabled (intermixing), whereby a subsurface alloy forms in which Ta accumulates in the second Pt surface layer, while bulk diffusion is not observed. Both adsorbed Ta adatoms on the Pt surface, and substitutional Ta in the Pt subsurface layer strongly interact with Pt and display a positive charge (+2). Hence, Ta atoms oxidize in contact with Pt, which is reflected in a substantial CLS to higher binding energy (BE) (1.5–2.0 eV) comparable to that of undercoordinated Ta atoms of exposed Ta single-crystal surfaces.<sup>28</sup> This alloy can be oxidized in an oxygen atmosphere at elevated temperatures (900 K) to yield Ta-oxide islands (surface segregation). Ta-oxide islands and Ta–Pt alloy can be converted into one another by reductive intermixing and oxidative segregation of Ta.

In this article, we extend our study to size-selected supported Ta clusters, elucidating the stability of Ta<sub>*n*</sub> (*n* = 4, 5, 6, 8, 13) clusters on Pt(111) in the presence and absence of oxygen. Performing the HR-XPS measurements at 40 K restricts surface diffusion of the monitored Ta species and strongly reduces phonon-induced broadening of the core-level spectra. Coupled with the overall energy resolution (better than 50 meV), fundamental spectroscopic insights can be derived from the rich and complex Ta 4f spectra of supported Ta clusters. Furthermore, we can exploit oxidation through atomic oxygen. This method<sup>29</sup> ensures access to the highest oxidation state, which is of particular interest concerning the observed BEs for Ta-oxides reported in our previous work,<sup>27</sup> indicating an

oxidation state lower than Ta<sup>5+</sup>. Based on the combined HR-XPS and DFT results, an oxidation mechanism is proposed for Ta<sub>8</sub> clusters, providing structural and electronic insights into the formation of the resulting oxidation product common to all cluster sizes. The product of oxidation is further characterized by room-temperature STM measurements in which the comparability to the low-temperature core-level measurements is reasoned in terms of the high similarity of the XPS spectra after oxidation at both low and high temperatures. Finally, and conclusive with our findings on Ta single atoms, a reductive intermixing of the oxidation product and subsequent oxidative surface segregation is observed.

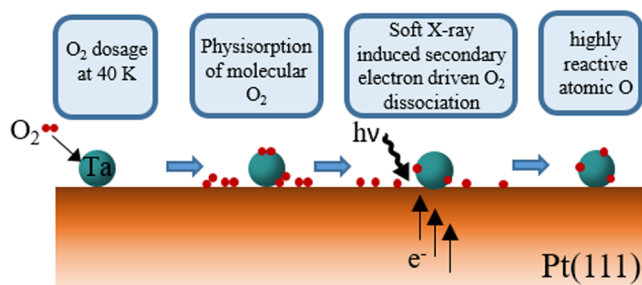
## EXPERIMENTAL SECTION

**Sample Preparation.** All XPS experiments were performed at the SuperESCA beamline of the Elettra synchrotron radiation facility (Trieste, Italy) in a UHV chamber at a 1 × 10<sup>−10</sup> mbar base pressure. A Pt(111) single crystal was prepared by subsequent cycles of sputtering with Ar<sup>+</sup>-ions (*E*<sub>kin</sub> = 1.5 keV, *p*<sub>Ar</sub> = 3.6 × 10<sup>−6</sup> mbar, *I* = 10 μA, 300 K, 20 min) followed by annealing >900 K for 1 min to ensure surface reconstruction (*T*<sub>Hüttig</sub> = 608 K).<sup>30</sup> Carbon impurities were removed by annealing in an O<sub>2</sub> atmosphere (*p*<sub>O<sub>2</sub></sub> = 1 × 10<sup>−7</sup> mbar, 870 K, 5 min), followed by the removal of residual adsorbed oxygen in H<sub>2</sub> atmosphere (*p*<sub>H<sub>2</sub></sub> = 5 × 10<sup>−8</sup> mbar, 670 K, 5 min). The process was repeated until no more contaminations were identified in the O 1s, C 1s, Ta 4f, and Pt 4f regions. The heating of the sample is enabled by electron bombardment from resistively heated W-filaments placed close to the back of the sample. The cluster source setup ENAC (Exact Number of Atoms in each Cluster) developed by the Nanoscale Materials Laboratory of the Department of Physics of the University of Trieste and Elettra Sincrotrone Trieste, which has proven very successful for similar cluster studies,<sup>29,31–34</sup> was employed for cluster generation. The source and the ion optics for charge and mass selection are based on the design of Heiz et al. and are described in detail elsewhere.<sup>35</sup> For the purpose of the following experiments, the cluster source was connected to the preparation chamber of the SuperESCA beamline, where sample cleaning and cluster deposition were performed. A valve gate separates the latter chamber from the analysis chamber, where the HR-XPS measurements of the as-deposited Ta species are conducted. Thus, sample preparation and analysis are carried out in situ. Furthermore, the chamber is connected to a gas manifold, enabling the exposure of the sample to reactive gases while monitoring changes in the photoemission spectra. Atomically precise Ta clusters of distinct sizes were deposited on a Pt(111) crystal at 40 K. The following cluster coverages per ML of Pt (ML<sub>Pt</sub>) were used: Θ(Ta<sub>4</sub>) = 0.10% clusters/ML<sub>Pt</sub>, Θ(Ta<sub>5</sub>) = 0.13% clusters/ML<sub>Pt</sub>, Θ(Ta<sub>6</sub>) = 0.11% clusters/ML<sub>Pt</sub>, Θ(Ta<sub>8</sub>) = 0.10% clusters/ML<sub>Pt</sub>, Θ(Ta<sub>13</sub>) = 0.06% clusters/ML<sub>Pt</sub> with respect to the atom density of the Pt(111) surface which corresponds to 1.5 × 10<sup>15</sup> atoms/cm<sup>2</sup>. Before deposition, the substrate was flash annealed up to 500 K to remove adsorbates accumulated from the chamber residual gas background. Ta atoms (Θ(Ta<sub>1</sub>) ≈ 2% atoms/ML<sub>Pt</sub>) were evaporated on the Pt(111) surface from a Ta wire (thickness: 0.125 mm, purity: 99.9%, Goodfellow) as described elsewhere.<sup>27</sup>

All STM measurements were performed in a different UHV chamber at 1 × 10<sup>−10</sup> mbar base pressure at the Technical

University of Munich. The Pt(111) single crystal was prepared by cycles of sputtering with Ar<sup>+</sup>-ions ( $E_{\text{kin}} = 1 \text{ keV}$ ,  $p_{\text{Ar}} = 4.0 \times 10^{-5} \text{ mbar}$ ,  $I = 10 \text{ }\mu\text{A}$ , 300 K, 20 min), followed by annealing >900 K for 5 min. Subsequently, the crystal was annealed at 720 K in an O<sub>2</sub> background ( $p_{\text{O}_2} = 5 \times 10^{-7} \text{ mbar}$ , 10 min) and ramped to 1000 K in UHV. Similarly to the other experiments, the Ta clusters were deposited with a cluster generation setup based on the design of Heiz et al.<sup>35</sup> under soft-landing conditions at room temperature. A Ta<sub>4</sub> coverage of 0.13% clusters/ML<sub>Pt</sub> and a Ta<sub>8</sub> coverage of 0.26% clusters/ML<sub>Pt</sub> have been obtained.

**Oxidation.** The oxidation of evaporated Ta atoms and deposited Ta clusters in the HR-XPS measurements at 40 K was performed according to a previously reported procedure<sup>29,31</sup> schematically depicted in Figure 1. Oxygen was



**Figure 1.** Scheme of the low-temperature oxidation procedure using physisorbed molecular O<sub>2</sub>. The soft X-ray radiation causes the release of secondary electrons that lead to the dissociation of O<sub>2</sub>, which forms highly reactive atomic oxygen.

predosed at 40 K (0.1–10 L), yielding a physisorbed layer of molecular O<sub>2</sub>. Irradiation of the sample with soft X-rays leads to the release of secondary electrons from the substrate that induce dissociation of the physisorbed O<sub>2</sub> molecules to form highly reactive atomic oxygen, which permits reaching the highest accessible oxidation state of the adsorbed Ta particles. Oxidation at elevated temperature (900 K) was performed in a background of molecular O<sub>2</sub> ( $p_{\text{O}_2} = 1 \times 10^{-6} \text{ mbar}$ ) instead. The transformation of the Ta-oxide islands into the Ta–Pt alloy was achieved by UHV annealing at 900 K until no further changes in the photoemission spectra were distinguished (<10 min). The cluster oxidation in the STM experiments was performed in a background of molecular O<sub>2</sub> ( $p_{\text{O}_2} = 1.2 \times 10^{-7} \text{ mbar}$ , 15 min) at room temperature. Thereafter, the sample was annealed at 600 K for 15 min in UHV.

**HR-XPS Measurements.** In the photon energy range applied here for the Ta 4f and Pt 4f spectra, a resolution better than 50 meV can be achieved.<sup>36</sup> Photoelectrons were collected and filtered using a Phoibos 150 mm mean-radius hemispherical electron energy analyzer (SPECS, Germany) in fixed analyzer transmission mode and detected with a 1D delay line detector. The following set of orbitals was probed at the corresponding photon energies indicated in brackets: O 1s ( $h\nu = 650 \text{ eV}$ ), C 1s ( $h\nu = 400 \text{ eV}$ ), Pt 4f ( $h\nu = 220 \text{ eV}$ ), and Ta 4f ( $h\nu = 150 \text{ eV}$ ). All acquired spectra were measured in normal emission geometry. The analyzer was operated in fixed analyzer transmission at a pass energy of 5 eV. Spectra were normalized to the photon flux of the incident beam and the number of scans. The core–electron binding energies of the experimental spectra were referenced to the Fermi energy, measured subsequently and under the same conditions. To

compare the computational core–electron BEs to the experimental ones, they have been aligned using the Ta–Pt subsurface alloy formed at 900 K in UHV as an internal experimental reference for Ta (IER<sub>Ta</sub>). This alloy species appears in the experiment with a well-defined peak at a BE of 23.59 eV, observed under the same measurement conditions as all the other Ta species discussed in this article.<sup>27</sup> The calculations are then referenced by taking the calculated BE of the defect-free Ta(110) slab (19.35 eV) and the Ta 4f CLS of the calculated subsurface alloy (+1.47 eV) with respect to the slab for realigning all calculated Ta 4f BEs. This leads to a realigned calculated BE for the Ta(110) photoemission of 22.12 eV, which is in agreement with experimental literature values.<sup>28</sup> Similarly, the calculated Pt 4f BEs of surface atoms in Pt(111) can be calibrated according to the experimental value of 70.52 eV corresponding to the surface component in Pt(111).<sup>37</sup>

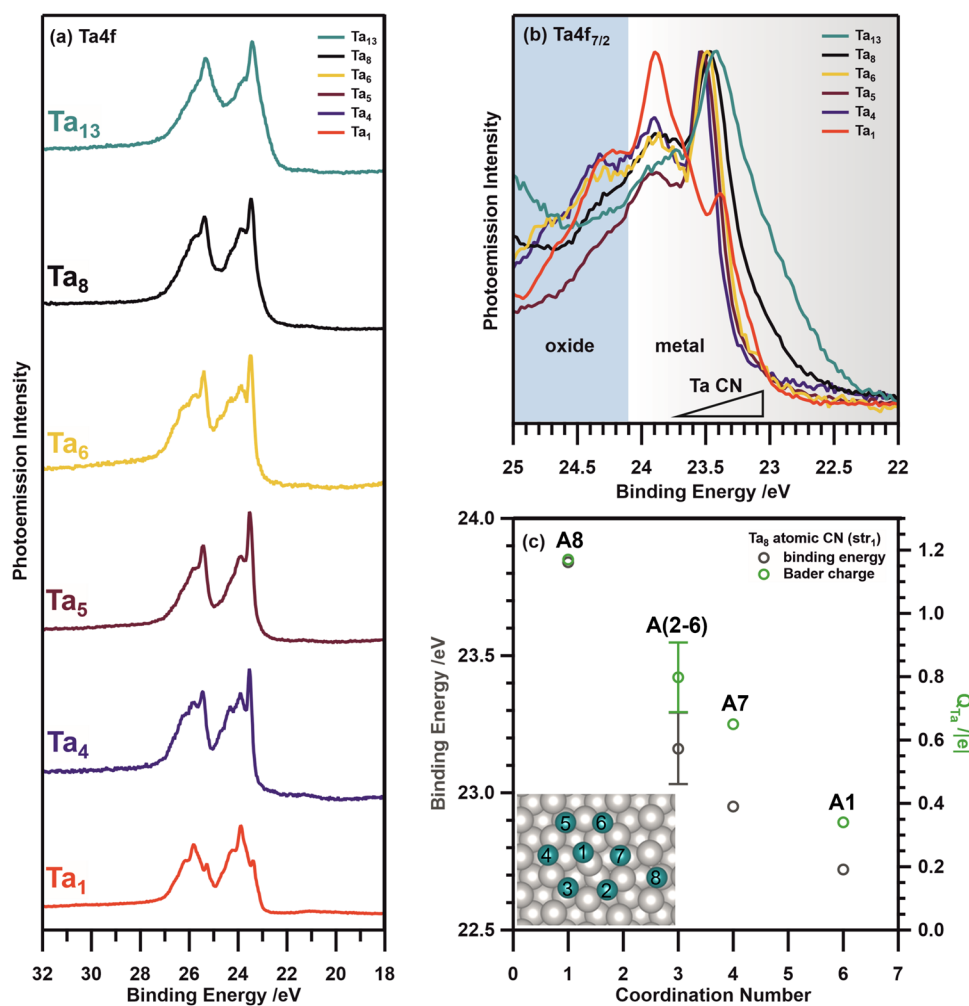
**STM Measurements.** The STM images were acquired in constant current mode with a commercial Scienta Omicron VT-AFM instrument, using homemade etched W tips. The images were processed with Gwyddion using plane subtraction and row-by-row alignment tools for background correction.<sup>38</sup> The height distribution of the particles was determined using a home-written Igor Pro Wavemetrics routine by detecting the particles via an intensity threshold, drawing a profile through the cluster maximum, and determining the cluster height with respect to the median background of the image.

**Computational Details.** All calculations were performed with the code VASP6.<sup>39,40</sup> The interaction between the core electrons and the nuclei is described with the PAW formalism.<sup>41,42</sup> The plane-wave basis set was expanded up to a kinetic energy of 400 eV. Truncation criteria of  $10^{-5} \text{ eV}$  for the electronic loop and 0.01 eV/Å for the ionic loop have been set. A PBE exchange–correlation functional was adopted.<sup>43</sup> The core-level energies have been calculated within the initial state approximation and realigned to the experimental BEs using the pure Ta(110) surface as a common reference. This approach is often found to yield a reasonable agreement with experimental findings for both metal adlayers and bulk alloys<sup>44</sup> and is adequate in particular when oxidation state changes lead to BE differences up to several eV. A  $5 \times 5 \text{ Pt(111)}$  cell has been adopted to accommodate the Ta<sub>4</sub> clusters, while an  $8 \times 8$  cell has been adopted for Ta<sub>8</sub>. The reciprocal space has been sampled with a Monkhorst net of  $4 \times 4 \times 1$  and  $2 \times 2 \times 1$  *K*-points, respectively. Pt(111) has been modeled with a five-layer slab, where the atoms from the bottom layer are frozen in their bulk positions while all other atoms are free to relax. An empty space of at least 20 Å is included in all supercells to avoid spurious interactions with the replica. The adsorption energy per Ta atom of Ta<sub>*n*</sub> species on Pt(111) has been calculated as follows

$$E_{\text{ads}} = \{E(\text{Ta}_n)/\text{Pt(111)} - [E(\text{Pt(111)}) + E(\text{Ta}_n)]\}/n$$

where Ta<sub>*n*</sub>/Pt(111) refers to a Ta<sub>*n*</sub> species bound to Pt, Pt(111) is the clean Pt surface, and Ta<sub>*n*</sub> is the gas-phase tantalum cluster in its most stable geometry and electronic ground state.

The Ta-oxide islands were modeled recurring to a two-dimensional tantalum oxide slab supported on Pt(111). The D3 Grimme parameterization and the Becke–Johnson damping function were adopted to account for the metal-oxide noncovalent interactions.<sup>45,46</sup> We focused on the β-Ta<sub>2</sub>O<sub>3</sub> phase, whose most stable surface is the (100), as assessed in



**Figure 2.** (a) Stacked overview of the Ta 4f XPS spectra ( $h\nu = 150$  eV) of Ta atoms directly after evaporation and Ta<sub>n</sub> clusters ( $n = 4, 5, 6, 8, 13$ ) directly after deposition on Pt(111) at 40 K. From bottom to top are displayed: Ta adatoms (red), Ta<sub>4</sub> (dark blue), Ta<sub>5</sub> (brown), Ta<sub>6</sub> (yellow), Ta<sub>8</sub> (black) and Ta<sub>13</sub> (turquoise). All Ta peaks are found to be considerably shifted toward higher BE with respect to metallic Ta<sup>0</sup> (21.64 eV).<sup>28,48,49</sup> (b) Cut-out of the Ta 4f<sub>7/2</sub> region for Ta<sub>n</sub> clusters ( $n = 4, 5, 6, 8, 13$ ) after the spectra are normalized to their maximum. The shadings indicate two main BE regions related to metallic (gray gradient) and oxidic (blue) Ta contributions. The gradient over the metallic region reflects the change in BE as a function of coordination within the Ta clusters as discussed in the main text. (c) Simulated BEs (green) and Bader charges (gray) of the Ta atoms (A1–A8) in a Ta<sub>8</sub> cluster deposited on Pt(111) (see inset) depending on the CN of the atoms. Atoms A(2–6) represent an average of atoms with the same CN.

a previous computational work.<sup>47</sup> However, due to the very poor lattice match between Ta<sub>2</sub>O<sub>5</sub>(100) and Pt(111), the interface with Pt(111) was simulated by cutting the oxide along another, less stable, low-index surface, namely (001). A model with a minor lattice strain (+1.7%, +1.3%) released on the oxide film was then constructed with a (3 × 4) Ta<sub>2</sub>O<sub>5</sub>(001) on a 4 × √3 Pt(111) coincidence lattice.

## RESULTS AND DISCUSSION

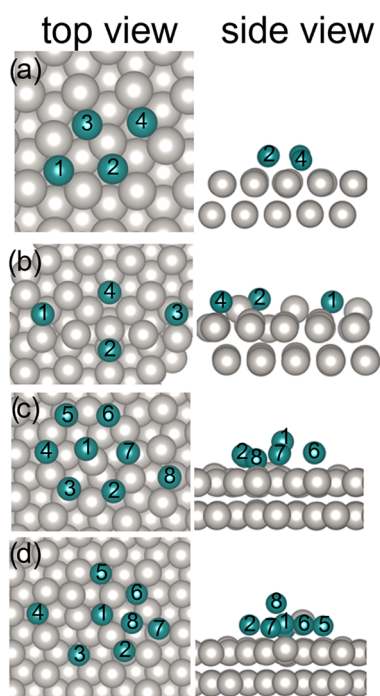
**Ta Clusters.** Ta<sub>n</sub> clusters ( $n = 4, 5, 6, 8, 13$ ) were deposited in UHV, under soft-landing conditions ( $E_{\text{kin}} < 1$  eV/atom), on a Pt(111) single crystal at 40 K using the laser ablation cluster generation setup ENAC.<sup>29</sup> The corresponding 4f core-level spectra of the as-deposited Ta<sub>n</sub> cluster sizes are shown in Figure 2a as a stacked overview and in Figure 2b as normalized to their maximum intensity. Duplets with a characteristic spin–orbit coupling of 1.91 eV and the 4:3 spin multiplicity intensity ratio (corresponding to Ta 4f<sub>7/2</sub> and Ta 4f<sub>5/2</sub>) are observed. The Ta 4f<sub>7/2</sub> peaks appear in the range 22–25 eV, where they can overlap with O 2s photoemission signals. All investigated Ta

4f emissions show a considerable shift to higher BEs with respect to emissions commonly observed for metallic Ta<sup>0</sup> bulk<sup>28,48,49</sup> (21.64 eV) and surfaces, e.g., Ta(100)<sup>48</sup> (22.39 eV), Ta(110)<sup>28</sup> (21.97 eV) and Ta(111)<sup>49</sup> (22.04 eV). This is similar to the CLSs found for Ta adatoms on the same support, where the interaction between the early transition metal Ta and the noble metal Pt leads to a formal oxidation of Ta by Pt. Ta adatoms are found to display an oxidation state of approximately +2, as suggested by simulations.<sup>27</sup>

Strong spectral similarities are found for all investigated Ta cluster sizes. Common photoemission regions are identified as indicated by the two shaded areas in Figure 2b. They can be attributed to contributions from metallic Ta clusters (gray gradient), ranging from 22.0 to 24.1 eV and oxidic Ta species (blue), with a Ta 4f<sub>7/2</sub> BE > 24.1 eV. Overall, the photoemission peaks from metallic Ta clusters found below 23.6 eV are most pronounced and exhibit a size-specific trend: For smaller clusters, e.g., Ta<sub>4</sub>, a relatively sharp photoemission onset is observed at ~23 eV. In contrast, this onset gradually evolves for larger clusters (investigated up to Ta<sub>13</sub>), as

observed in the contributions at the low BE end of the spectra. These photoemission features are in good agreement with those observed for Ta agglomerates (23.4 eV), which have been previously identified after deposition of Ta atoms. The formation of agglomerates was assigned to the transient atom mobility due to residual  $E_{\text{kin}}$  of  $\sim 0.28$  eV/atom upon evaporation.<sup>27</sup> Comparison of the Ta cluster with the Ta adatom (red) spectra indicates a contribution similar to that observed for Ta adatoms (23.6–24.1 eV) and is present in the spectra of all cluster sizes. To further investigate the  $4f_{7/2}$  spectral trend observed for different Ta cluster sizes and the origin of contributions similar to Ta adatoms emissions to the cluster spectra, insights into cluster geometry and stability are required. Therefore, the structure and stability of Ta<sub>4</sub> and Ta<sub>8</sub> on Pt(111) have been investigated by DFT and STM.

For Ta<sub>4</sub> clusters, a gas-phase simulation returns a tetrahedral shape, see Figure S1 left, and a closed-shell ground state configuration.<sup>50</sup> The adsorption energy ( $E_{\text{ads}}$ ) per Ta atom, however, reveals that a flat isomer,  $\text{str}_1(\text{Ta}_4)$ , is favored ( $E_{\text{ads}} = -3.25$  eV/atom) on the Pt(111) support, Figure 3a, and



**Figure 3.** Computed Ta<sub>4</sub> and Ta<sub>8</sub> structures on Pt(111). Top views (left column) and the respective side views (right column) of the structures  $\text{str}_1(\text{Ta}_4)$  (a),  $\text{str}_2(\text{Ta}_4)$  (b),  $\text{str}_1(\text{Ta}_8)$  (c), and  $\text{str}_2(\text{Ta}_8)$  (d).

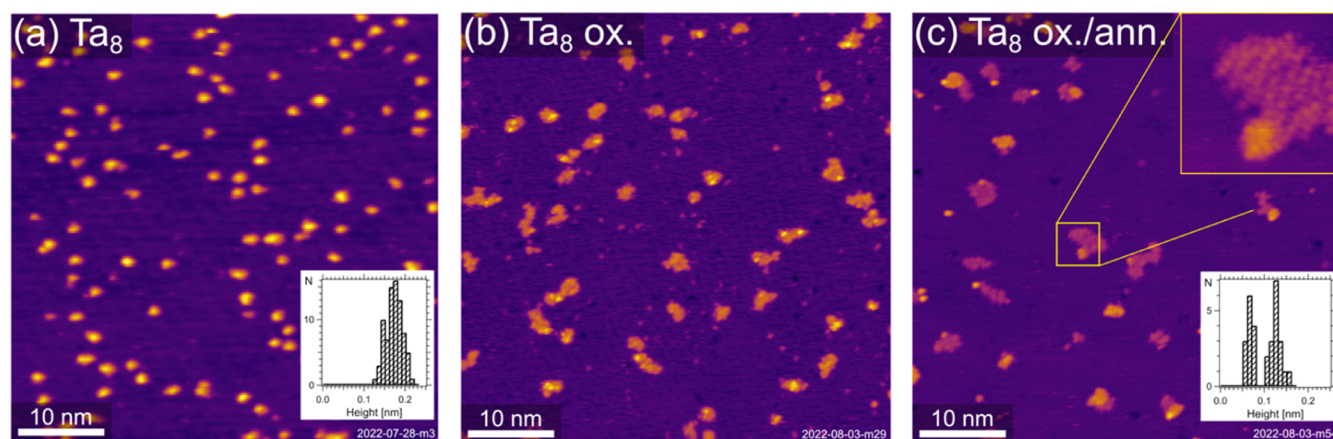
tetrahedral structures are much less stable. Two of the Ta atoms (A1, A4) of the flat tetramer are in good registry with the underlying hollow sites and display a core–electron BE of 23.43 eV, see Table S1, slightly smaller than that of Ta adatoms on the surface. The remaining two atoms (A2, A3) display a poorer registry and a smaller BE of 22.93 eV. Further energy is gained if one Ta atom takes over for Pt in a lattice site:  $\text{str}_2(\text{Ta}_4)$  has an  $E_{\text{ads}}$  of  $-3.96$  eV/atom, see Figure 3b. However, given the low temperature of the experiments this rearrangement will likely be experimentally inaccessible. The Ta–Pt substitution that contributes to the overall stabilization is supposedly hindered by an energy barrier that is hard to account for in the calculations. STM measurements on the Ta<sub>4</sub>

cluster on Pt(111) (0.13% clusters/ML<sub>Pt</sub>, 293 K), see Figure S2, reveal that, indeed, a flat geometry is adopted on Pt(111), and the height profile indicates an apparent height for Ta<sub>4</sub> clusters of  $\sim 0.12$  nm. While intact, monodisperse Ta<sub>4</sub> clusters dominate, some particles accumulated at step edges and sintered ones on the terraces are also observed. Furthermore, some fragments of unspecified size can be found. In this context, it is interesting to point out that the  $E_{\text{ads}}$  for Ta<sub>4</sub> is directly comparable to the case of Ta adatoms. A crosscheck immediately reveals that a fully dispersed arrangement of atomic Ta species on the surface ( $E_{\text{ads}} = -3.84$  eV/atom, Table S1) is of comparable stability to that of the flat Ta<sub>4</sub> structure. Thus, in the case of Ta<sub>4</sub>, the thermodynamic drive toward dispersion seems to be kinetically accessible at room temperature, as indicated by the fragments observed in the STM images. Nevertheless, at 40 K and under irradiation, no evidence for X-ray beam damage-induced fragmentation is found, neither for Ta<sub>4</sub> clusters nor any other investigated cluster size, as shown in Figure S3. Hence, the origin of atomic contributions in cluster XPS spectra at 40 K cannot be explained by probe-induced fragmentation and indicates a different origin.

In the gas-phase, Ta<sub>8</sub> assumes the shape of a capped octahedron;<sup>50</sup> see Figure S1 right. On the Pt(111) surface, however, our calculations show that this arrangement is not stable upon adsorption, and the Ta<sub>8</sub> cluster again, similar to Ta<sub>4</sub>, tends to spread on the surface, maximizing the interaction with Pt. In particular, two minima with different structures have been identified, namely  $\text{str}_1(\text{Ta}_8)$  and  $\text{str}_2(\text{Ta}_8)$ , which are displayed in Figure 3c,d as top and side views, respectively.  $\text{str}_1(\text{Ta}_8)$  appears as an almost flat structure, with a maximum height of 0.29 nm with respect to the Pt apical plane. Six Ta atoms (A2–A7) occupy bridge Pt–Pt sites in an almost hexagonal arrangement; one Ta atom is hosted in a Pt<sub>3</sub>-hollow side in the middle of the hexagon (A1), and the last one is adsorbed aside on a hollow position (A8). The  $E_{\text{ads}} = -2.13$  eV/atom, thus indicates a substantial thermodynamic gain related to the adsorption and almost complete flattening of the Ta<sub>8</sub> cluster on Pt(111). The calculated Ta 4f BEs span from 22.72–23.84 eV, see Table S2. The corresponding Bader charges and interatomic Ta–Ta distances are reported in Tables S2 and S3, respectively.

Similarly to the case of Ta<sub>4</sub>, an even more stable arrangement,  $\text{str}_2(\text{Ta}_8)$ , could be identified and is shown in Figure 3d: Interestingly, during the structure relaxation, spontaneous incorporation of a Ta atom (A4) into Pt(111) took place, and the expelled Pt atom remained aggregated to the Ta cluster. In this case, the adsorption energy is  $-2.68$  eV/atom.  $\text{str}_2(\text{Ta}_8)$ , thus, is in principle more stable than  $\text{str}_1(\text{Ta}_8)$ . In our previous work,<sup>27</sup> we showed how the Ta–Pt substitution is thermodynamically favorable in surface sites, and an even stronger stabilization is envisaged if the Ta impurity migrates to a bulk Pt site, in agreement with a previous computational work.<sup>51</sup> However, like for Ta<sub>4</sub>, Ta–Pt substitution to form  $\text{str}_1(\text{Ta}_8)$  is probably hindered at 40 K. The alloying with Pt induces a general increase in the calculated Ta 4f BEs, see Table S1. Notably, both structures discussed here display a quasi-flat arrangement of the Ta atoms, compatible with apparent heights typical of adsorbed monolayer structures.

The STM measurements of Ta<sub>8</sub> clusters, deposited and recorded at room temperature, on Pt(111) in Figure 4a display monodisperse particles with a mostly flat two-dimensional



**Figure 4.** STM images of Ta<sub>8</sub> clusters on Pt(111) illustrate the morphology of deposited clusters and their changes upon oxidation and/or annealing. The clusters were deposited and measured at room temperature (0.26% clusters/ML<sub>Pt</sub>). (a) Display of the as-deposited pristine Ta<sub>8</sub> clusters. The height histogram (inset) reveals the monodisperse formation of clusters with an apparent height of 0.18 nm. (b) After oxidation at room temperature ( $p_{\text{O}_2} = 1.2 \times 10^{-7}$  mbar, 15 min), the clusters have considerably sintered to form larger, flatter oxidic species (0.1% species/ML<sub>Pt</sub>), on which occasionally small, higher features reside. (c) After further annealing ( $p_{\text{O}_2} = 1.2 \times 10^{-7}$  mbar, 600 K, 15 min), which also leads to the partial desorption of oxygen from the Pt surface, the image resolution slightly improves (see, e.g., the enlarged inset). The dominant flat terraces now have an apparent height of 0.07 nm, while small second-layer areas appear at a height of 0.13 nm (see histogram inset). Imaging parameters for all images:  $U_{\text{bias}} = 1.5$  V,  $I_t = 0.3$  nA,  $50 \times 50$  nm<sup>2</sup>.

(2D) geometry. The height histogram sampled over a terrace (inset of Figure 4a) indicates an apparent cluster height of  $\sim 0.18$  nm, which is smaller than an atomic step height of a Pt(111) surface (0.226 nm).<sup>52</sup> This lower apparent height of the Ta clusters could be related to different electron densities of states of the two involved metals that considerably affect the observed heights in the STM topography of heteroepitactic contacts. The clusters are found intact (over several hours) with no conclusive indication of the presence of Ta fragments (<Ta<sub>8</sub>) or Ta adatoms on the Pt surface, even at room temperature. Additionally, under X-ray irradiation (150 eV), negligible changes are observed over a series of five XPS spectra ( $\sim 8$  min), see Figure S3. Hence, considering the sequence of spectra and STM, atomically precise Ta<sub>8</sub> clusters are stable upon deposition at  $T < 293$  K and under irradiation at 40 K.

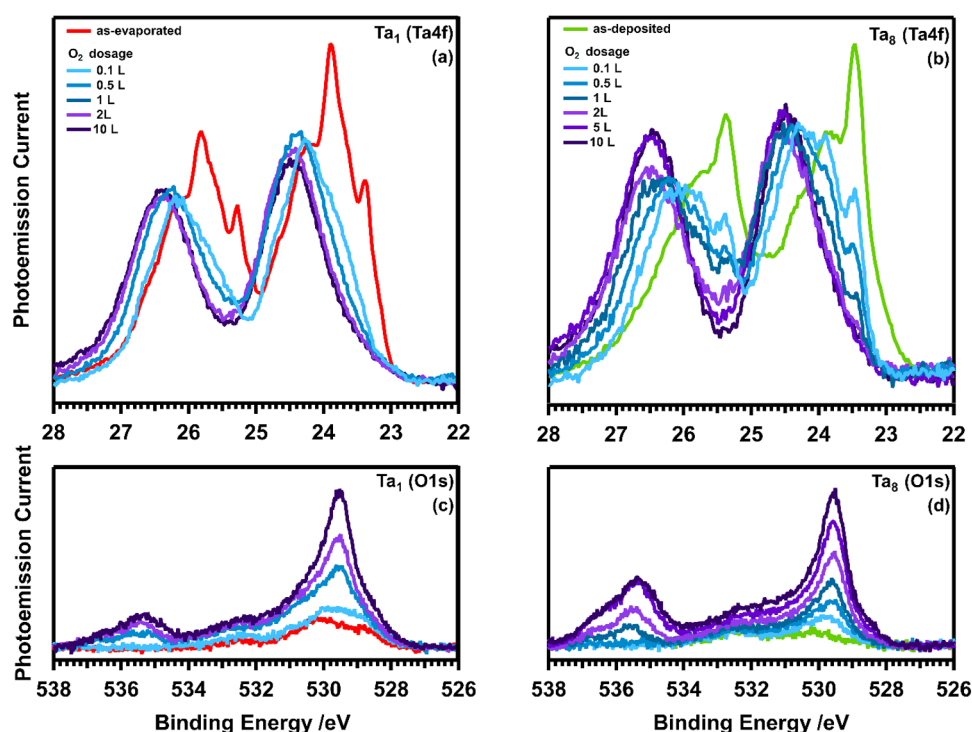
Overall, considering the flat structure of Ta<sub>4</sub> and Ta<sub>8</sub>, as well as the common trend of a preferred 2D geometry for metal clusters on metal supports<sup>53</sup> to minimize surface free energy, such as observed for Ag<sub>19</sub> on Pt(111)<sup>54</sup> or Pd<sub>19</sub> on Rh(111),<sup>55</sup> a flat geometry could be assumed for other Ta cluster sizes too. The tendency of Ta clusters of this size to flatten and spread on the Pt surface is due to the stronger Pt–Ta interactions with respect to the cohesive energy of the Ta clusters.

In this regard, the differences in photoemission of Ta clusters of distinct size in Figure 2a,b could be interpreted as follows:

Considering 2D structures of Ta<sub>n</sub> clusters on Pt(111), their arrangement can be understood as a flat shell structure, where the central Ta atoms are more coordinated than the peripheral atoms, analogously to the trend observed from surface Ta atoms of reduced coordination on open surfaces. Here, an increase in BE is observed for Ta atoms of reduced coordination.<sup>28,56</sup> The Ta–Ta residual coordination is proved by their calculated reciprocal distances being, in many cases, comparable to what observed for gas-phase Ta<sub>8</sub> (Tables S3 and S4), with a general elongation of 0.02 nm to better fit with the underlying Pt(111) surface. In the following, str<sub>1</sub>(Ta<sub>8</sub>) is

employed to further elucidate the XPS spectra in terms of distinct in-plane coordination number (CN) for Ta atoms in a Ta cluster. Based on the bond length found for cationic Ta clusters in the gas-phase (0.24–0.29 nm<sup>57</sup>)—we defined a truncation criterion of 0.35 nm for Pt(111) supported Ta clusters. Hence, Ta atoms displaying a Ta–Ta bond  $< 0.35$  nm, see Table S3, are considered as coordinated to the neighboring Ta atom. Consequently, a CN can be assigned to the individual Ta atoms (A1–A8) and they can be divided into four sets of in plane coordinated atoms: A1 (CN = 6), A7 (CN = 4), A8 (CN = 1) and the group of equivalent coordination A2–A6 (CN = 3), see Figure 2c.

In this sense, the contribution to the lowest BE (22.72 eV) arises from A1—the central tantalum atom of the cluster, see inset of Figure 2c. Vice versa, A8, which is adsorbed in a hollow site and solely coordinated to A7, displays a larger calculated BE (23.84 eV), a value comparable with what was reported for isolated Ta adatoms on Pt(111).<sup>27</sup> All Ta atoms composing the hexagonal arrangement (A2–A7) around the central Ta atom exhibit intermediate BEs (23–23.4 eV), see Table S2. The dispersion of values in this range is most likely due to the effect of Ta–Ta interactions and different adsorption configurations with respect to the Pt(111) surface. Noteworthy, due to the additional coordination to A8, A7 displays the lowest BE of the hexagonal arrangement. Consequently, emissions from less coordinated Ta atoms are found to shift to higher BE and vice versa. Therefore, central atoms in a Ta<sub>13</sub> cluster express a more bulk-like behavior than, for example, atoms in a Ta<sub>4</sub> cluster, as they are neighbored by more Ta atoms and of higher coordination. Hence, the photoemission gradually broadens at the lower BE onset with increasing cluster size.<sup>31,34</sup> Therefore, we can attribute the features in the Ta<sub>n</sub> spectra in Figure 2b to atoms with different CN in the clusters, whose density depends on the clusters size. Features at lower BE (less than 23.2 eV) can be attributed to atoms with higher CN ( $> 4$ ), which are more abundant for larger clusters, the main peak from the spectra in the range 23.2–23.6 eV is associated with atoms with intermediate CN



**Figure 5.** Sequence of XPS spectra showing the evolution of the Ta 4f ( $h\nu = 150$  eV) emissions, where the oxidation process sets in upon irradiation. (a) As-evaporated Ta atoms and (b) as-deposited  $\text{Ta}_8$  clusters on Pt(111) upon subsequent exposure to 0.1–10 L  $\text{O}_2$  at 40 K. The O 2s measurements for 0.1–10 L on Pt(111), see Figure S5, were subtracted from all oxidized Ta 4f spectra. The bottom spectra (c, d) display the corresponding evolution of the O 1s ( $h\nu = 650$  eV) region.

(2–4) which are expected to be abundant in all the clusters due to their planar geometry, see the case of  $\text{Ta}_8$  and  $\text{Ta}_4$ . Finally, the spectral region in the range 23.6–24.1 eV is attributed to atoms with the lowest CN (<2). This component is the highest in the case of  $\text{Ta}_1$ , as expected from our interpretation, but is not zero for the larger clusters such as  $\text{Ta}_8$  where it is still possible to have atoms with low CN in the outer regions of the clusters, see Figure 2c. On the contrary, the small components in the region associated with intermediate and high CN in the  $\text{Ta}_1$  spectra are due to small agglomerates upon deposition of the adatoms.

Interestingly, the Bader charge expresses a similar trend as the BE with respect to the CN, see Figure 2c, and correlates with the Ta–Pt distance of the respective Ta atom, see Table S2. A1 is the most distant from the Pt surface and exhibits the lowest charge (+0.34|e|), while A8 is the closest to the Pt surface and possesses the highest charge (+1.17|e|). Consequently, an increase in both factors, the CN and Ta–Pt distance, contribute to decreasing the Bader charge and hence affect the oxidation state of the Ta cluster atoms. Furthermore, A8 presenting the highest charge while displaying the lowest CN is also in agreement with our findings concerning single Ta adatoms on Pt(111), where no possibility for in-plane coordination with Ta is given and an overall charge of +2 was identified.<sup>27</sup>

Among all of the Ta cluster sizes and the Ta adatoms, similar photoemission features are found on the high BE side (>24.1 eV) of the spectra. Considering the BE values of the Ta  $4f_{7/2}$  photoemission of Ta-oxides in the literature (22.5–27 eV<sup>6,16</sup>), we attribute all the features observed above 24.1 eV to Ta-oxides—this region is indicated in Figure 2b in blue. The observed spectral structure in the range 24.1–25 eV suggests contributions from multiple oxidic structures or configurations,

as inflection points can be observed. The distinct contributions to the Ta-oxide photoemission signals become more pronounced upon oxidation and have been discussed in ref 27. and in more detail in the following.

**Oxidation.** The different Ta cluster sizes and Ta adatoms were subsequently exposed to 0.1–10 L molecular  $\text{O}_2$  at 40 K (see Experimental Section). Upon oxidation emissions from both Ta 4f and O 2s orbitals can be identified in the region 22–28 eV. Although, the main features can be assigned to Ta 4f emissions, emissions from O 2s orbitals may also contribute to the measured spectra. This is especially true at higher  $\text{O}_2$  exposure ( $\geq 1$  L), as evident from O 2s spectra recorded under the same conditions, in which the clean Pt(111) crystal was exposed to 0.1–10 L  $\text{O}_2$ , see Figure S5. These O 2s spectra were subtracted from the Ta 4f spectra respective to the  $\text{O}_2$  dosage in order to account for the superposed emissions from chemisorbed molecular oxygen ( $\text{O}_{\text{chem}}$ ). The evolution of the respective Ta 4f and O 1s signals after deposition are displayed for Ta adatoms in Figure 5a,c and for  $\text{Ta}_8$  clusters in Figure 5b,d. Overall, in the Ta 4f emissions a gradual decline in emissions from both metallic Ta species of lower CN < 2 (BE 23.60–24.10 eV) and intermediate to higher CN > 2 (BE 23.60–22.00 eV), is observed, while features >24.10 eV originating from Ta-oxides become more pronounced. A similar progress is observed for all cluster sizes, as seen in Figure S4.

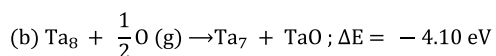
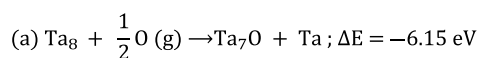
A low oxygen exposure of 0.1 L induces considerable changes in the Ta 4f and O 1s spectra of both Ta adatoms and  $\text{Ta}_8$  clusters. High CN features, i.e., components at BE < 23.2 eV in Figure 2c and associated with atoms with CN > 4, decrease considerably in intensity for  $\text{Ta}_8$  and even vanish in the case of  $\text{Ta}_1$ , where they derived from agglomerates formed upon evaporation. Thus, an already low exposure to oxygen

induces a rapid reaction, emphasizing the high oxygen affinity of Ta. The emissions related to atoms of low and intermediate CN (CN < 4 and BE = 23.2–24.1 eV) also decrease in intensity, but are sustained up to a higher exposure of 0.5 L in the Ta<sub>1</sub> spectrum and 1 L in the case of the Ta<sub>8</sub> spectrum. Thus, Ta<sub>8</sub> seems to exhibit lower susceptibility to oxidation than Ta<sub>1</sub>.

It is noteworthy that, for the Ta<sub>8</sub> clusters, the increase of photoemission signals of lower coordinated atoms could occur at the expense of higher coordinated ones. Thus, the sustained emission signal found at ~23.6–23.8 eV for Ta<sub>8</sub> could be reasoned by the formation of less coordinated Ta species upon cluster oxidation. In this scenario, an oxidative degradation of the Ta clusters is envisaged. Through abstraction—most likely initiated from Ta atoms of lower CN, i.e., from the periphery e.g., A8, atomic Ta species and concomitantly Ta cluster of reduced size would be yielded, resulting in Ta clusters comprising atoms of reduced coordination. Both Ta adatoms and Ta atoms of lower CN in a cluster can contribute to emissions found from 23.6 to 23.8 eV.

This consideration agrees with the observations from the STM image in Figure 4b, where Ta fragments are observed after oxidation next to Ta-oxide islands, which are discussed later. Gas-phase oxidation studies on cationic Ta clusters support this picture: Upon oxidation, not solely intact Ta-oxo clusters form, but degradation takes place as well, leading to the abstraction of fragments from the cluster that is observed in the size range Ta<sub>4</sub><sup>+</sup>–Ta<sub>12</sub><sup>+</sup>.<sup>58,59</sup> This oxidative cluster decomposition could account for the observed trend in emissions in parallel to the appearance of oxidized Ta species. In the gas-phase, for example, the oxidation of Ta<sub>9</sub><sup>+</sup> displays a considerably low activation energy (0.55 ± 0.03 kJ × mol<sup>-1</sup>) for the first degradation step.<sup>59</sup> Considering that highly reactive atomic O is formed upon irradiation (ΔH<sub>f</sub>(0 K) = 493.69 kJ/mol),<sup>60</sup> and thus the oxidation process is activated, and given the small barriers for oxidation of Ta clusters that are obtained from gas-phase experiments, it can be concluded that the oxidation of Ta clusters and atoms (ΔH<sub>f</sub>(297.15 K) = 192.46 kJ/mol)<sup>61</sup> under our experimental conditions are kinetically achievable. In the following Scheme 1, two

#### Scheme 1. Oxidation Mechanism with (a) Ta Atom and (b) TaO Abstraction



pathways, see also Figure S6, are proposed as tentative mechanisms for Ta cluster fragmentation upon reaction with atomic oxygen and have been treated computationally for the case of Ta<sub>8</sub> clusters.

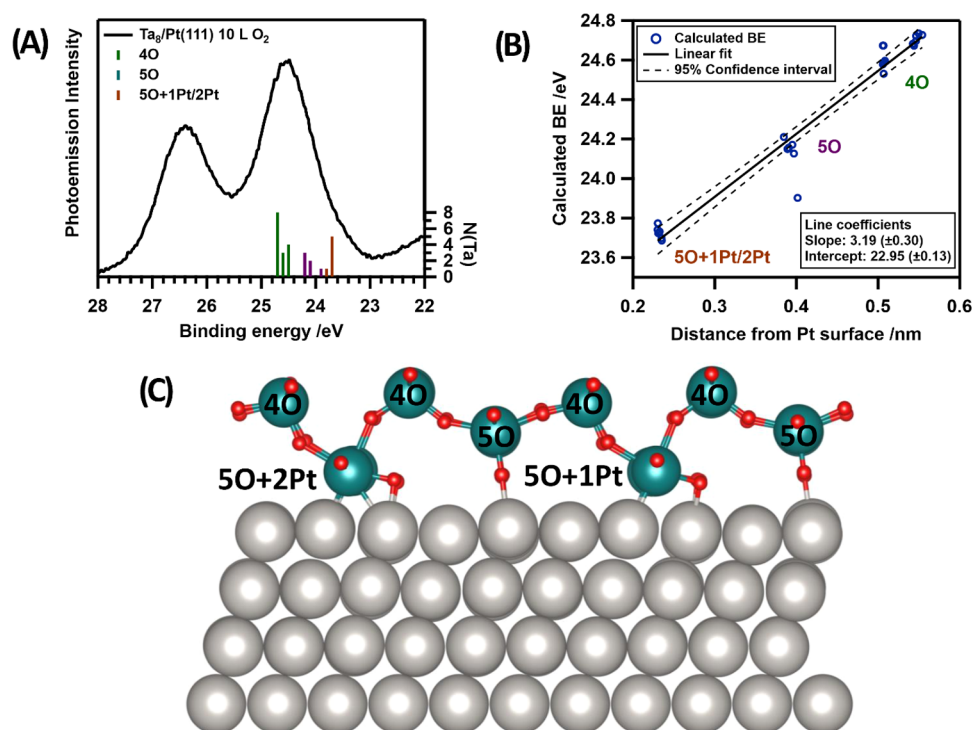
In the first mechanism, Scheme 1a, an oxygen atom is adsorbed on a Ta<sub>8</sub> cluster, mimicking the initial stage of the oxidation, and this results in a partially oxidized Ta<sub>7</sub>O adduct and an isolated Ta atom bound to Pt(111). In the second mechanism, Scheme 1b, a TaO species forms and leaves a smaller, unoxidized Ta<sub>7</sub> cluster. Simulation shows that both mechanisms are exothermic. However, the energetically more favorable mechanism is the abstraction of a Ta atom concomitant to oxidation of the remaining cluster of shrunk size as proposed in Scheme 1a. The resulting BE for Ta

adatoms (23.62 and 23.77 eV for hcp and fcp hollow sites<sup>27</sup>) agrees with the photoemission signals observed in the oxidation spectra in Figure 5b. For the oxidized Ta<sub>7</sub>O cluster, the oxygen atom is three-coordinated to two peripheral and central Ta atoms. We observe that the peripheral Ta atoms directly bound to oxygen display a larger BE (23.80 eV) with respect to what is observed for unoxidized clusters. The central Ta atom also undergoes a substantial increase in BE (23.04 eV) compared to metallic Ta<sub>8</sub>. The mechanism suggests that the features observed in the Ta<sub>8</sub> cluster spectra around 23.60 eV for higher O<sub>2</sub> exposures up to 1 L could be attributed to (i) Ta adatoms, (ii) Ta<sub>7</sub>O clusters (Table S5) as well as to (iii) further degradation products of subsequently fragmented Ta-oxo-clusters displaying a reduced overall coordination.

The presence of metallic Ta atoms after oxidation can be explained experimentally, as the oxidation method does not supply oxygen continuously. As a result, atomic oxygen may be unavailable for reaction either because it is consumed during irradiation before Ta atoms can form due to fragmentation or because it is bound and too distant from the produced atoms (immobilized at 40 K). Our findings suggest the oxidation to proceed along mechanism (a), but with no complete fragmentation of the Ta clusters. More details are discussed in the Supporting Information (SI) in Figure S6.

Upon further exposure to oxygen (≥2 L), the Ta 4f spectra contain no longer metallic Ta contributions but solely the Ta-oxide-related ones at 24.50 eV. The corresponding O 1s (hν = 650 eV) spectra, see Figures 5c,d and S5, indicate emissions from molecular O<sub>chem</sub> on Pt(111) at 529.60–530.54 eV, atomic oxygen at 529.83 eV and physisorbed molecular O<sub>2</sub> on Pt(111) at 534–538 eV.<sup>62</sup> Ta-oxide-related photoemission signals around ~530 eV are observed in the O 1s spectra of both, Ta adatoms and Ta<sub>8</sub> clusters, directly after deposition and in line with the corresponding Ta 4f signatures (>24.1 eV). Compared to literature, BEs of ~530 eV could be attributed to Ta<sub>2</sub>O<sub>5</sub> (530.2–530.4 eV),<sup>6,16</sup> which is also compatible with the high oxophilicity of Ta. With respect to Ta<sub>7/2</sub> in the Ta 4f spectrum, an oxidation state of Ta<sup>5+</sup>, is expected to appear at BEs > 26 eV,<sup>6,16</sup> which is clearly higher than the observed BEs (<25 eV). This might suggest that in our experiments the +5 state is not achieved, which would not be surprising based on several reported studies on the oxidation of metal clusters, which have shown that clusters are prone to oxidize in a 1:1 stoichiometry and hence to adopt a lower oxidation state compared to bulk materials.<sup>29,31–33</sup> However, our present DFT simulations indicate that indeed the oxidation state of Ta is +5 and that the shift toward lower BEs is caused by the interaction with the Pt support.

**Ta-Oxide Islands.** The observed XPS spectra of oxidized Ta adatoms and clusters at 40 K (see Figures S4 and S7) indicate the formation of a common oxidation product, independent of cluster size. To further elucidate, Ta<sub>8</sub>/Pt(111) samples were investigated by STM after oxidation and subsequent annealing, see Figure 4b,c. While the oxidation mechanism at low and high temperatures is not directly comparable, the XPS spectra after oxidation at 40 and 300 K display only small spectral differences, which are discussed later in this section; see Figure 7. In contrast to the stable and monodisperse clusters directly after deposition observed in Figure 4a, both measurements (b,c) indicate Ta agglomeration into oxidic islands. After oxidation, small Ta moieties are visible along the islands in Figure 4b. In contrast, these small moieties are no longer observed after subsequent annealing,



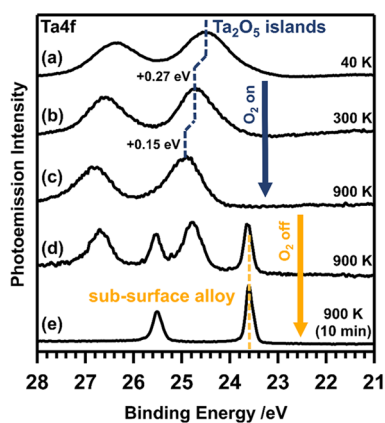
**Figure 6.** Computational model for Pt(111)-supported Ta<sub>2</sub>O<sub>5</sub> islands with distinct surface and interface contributions (4O, 5O, 5O+1Pt, and 5O+2Pt) of Ta with O and Pt. (A) The computed BEs are compared to the experimental BEs and found in good agreement. The vertical axis on the right N(Ta) indicates the relative count of the single contributions to the film. (B) The BEs of all contributions versus their height with respect to the Pt(111) surface are reported. A pronounced perturbation from the Pt surface on the Ta BEs is observed that is particularly strong for Ta atoms in contact with Pt. (C) The simulated structure exhibits surface and interface contributions. The surface layer is constituted of four-coordinate 4O (24.5–24.9 eV) Ta species. The interface layer is in contact with the Pt surface and composed of three nonequivalent five-coordinated Ta species, whereas 5O (23.9–24.3 eV), 5O+1Pt and 5O+2Pt (23.6–23.8 eV) contributions can be distinguished.

and the number of islands decreases while they increase in size, Figure 4c. The islands do not have a particular shape and appear ragged. They consist of ~50 atoms on average and are flat, similar to Ta clusters. The height profile reveals an apparent height of 0.07 nm with some second layer features of an apparent height of 0.13 nm. Hence, the Ta-oxide islands appear less high than the metallic Ta<sub>8</sub> clusters (~0.15 nm). It is intuitive to assign this island formation to the annealing process and cannot be excluded, however, as shown in the following, theoretical calculations suggest that for our samples island formations is also possible at 40 K. Experimentally, the formation of Ta-oxide islands at low temperatures, see Figure 7a, i.e., without thermal activation, can be reasoned as induced by the high oxophilicity of Ta. The nucleation into islands and lateral growth are known to often form the initial stage of oxidation, even at low temperatures.<sup>63</sup> The Ta-oxide islands observed in STM likely represent a product similar to the common oxidation product observed for oxidized clusters at 40 K.

We simulated oxidic Ta islands, adopting the Ta<sub>2</sub>O<sub>5</sub> stoichiometry in the model. The resulting film represents the thinnest possible continuous model and comprises a two-dimensional layer of Ta<sub>2</sub>O<sub>5</sub> building blocks, as shown in Figure 6. Upon coordination to Pt(111), the film exhibits a four-coordinate species (4O) at the surface and three nonequivalent five-coordinated species (5O, 5O+1Pt, 5O+2Pt) at the interface with Pt. The interface contributions differ in their local coordination regarding binding to the Pt(111) surface. Some of the Ta atoms bind to Pt through an O atom, while others adopt a direct coordination to the Pt surface. The

corresponding side view in Figure 6C reveals 3 different types of Ta coordination: 4-fold coordination to oxygen atoms (4O), 5-fold coordination to oxygen atoms (5O), and 5-fold coordination to O envisaging also one (5O+1Pt) or two (5O+2Pt) bonds to the Pt(111) surface. The BEs of these types of Ta coordination are reported in Figure 6A, which are in good agreement with observed XPS spectra recorded at 40 K for oxidized samples. This supports the presence of agglomerates already at low temperatures without thermal activation. The graph in Figure 6B demonstrates that their calculated BEs correlate linearly with the Ta atoms' height above to the Pt surface. This supports the strong perturbation exerted by the Pt surface on the Ta core levels, since the Ta atoms in contact to Pt display the smaller BEs (23.6–23.8 eV). A middle layer composed by 5O Ta species has intermediate BEs (23.9–24.3 eV). The top layer with four-coordinated Ta atoms shows the highest BEs (24.5–24.9 eV). The distance from Pt, thus, dominates over the influence of coordination and binding partners. This explains also the difficulty in associating a Ta oxidation state based on the observed Ta 4f BEs of the Ta-oxides.

Annealing the Ta-oxide islands formed at low temperatures in an oxygen atmosphere ( $p_{\text{O}_2} = 1 \times 10^{-6}$  mbar) to 300 K (+0.27 eV) and subsequently to 900 K (+0.15 eV) leads to a CLS and a narrowing of the Ta 4f photoemission signals, as shown in Figure 7a–c. The overall shift to higher BEs of +0.42 eV upon temperature increase points toward elongation of the Ta–O bonds with respect to the Ta<sub>2</sub>O<sub>5</sub> phase at 40 K.<sup>64</sup> The narrowing can be reasoned by converging to the minimum energy configuration, which becomes more accessible at higher



**Figure 7.** (a–c) Thermal evolution of the Ta 4f ( $h\nu = 150$  eV) emissions of the Ta-oxide islands from 40 to 900 K in an oxygen atmosphere ( $p(\text{O}_2) = 1 \times 10^{-6}$  mbar), followed by the conversion of the Ta-oxide islands into the Ta–Pt subsurface alloy under reductive UHV conditions at 900 K (c–e). Spectrum (d) represents a mixed-alloy state displaying emissions from both the Ta-oxide islands and the Ta–Pt subsurface alloy.

temperatures. When the oxygen supply is turned off at 900 K, the photoemission signals from Ta-oxide islands gradually decrease, while that of the Ta–Pt subsurface alloy (23.59 eV) increases, as shown for the intermediate mixed-island-alloy state shown in Figure 7d. Finally, after <10 min of further annealing in UHV, the Ta 4f spectrum in Figure 7e displays solely emissions from the subsurface alloy. As demonstrated in our previous study, this alloy can be converted back into Ta-oxide islands in the presence of oxygen.<sup>27</sup> Thus, two distinct chemical states of Ta at the Pt(111) surface can be interconverted by control of the environment (atmosphere and temperature), demonstrating the reversible duality of oxidative surface segregation and reductive subsurface intermixing.

## CONCLUSIONS

Size-selected Ta clusters soft-landed on a Pt(111) single crystal at 40 K were investigated using XPS and DFT. Concurrent photoemission signals from Ta clusters, adatoms and oxides have been spectroscopically observed upon deposition. The emissions of as-deposited Ta clusters display a size-specific trend. Emissions from Ta atoms with lower coordination, as in  $\text{Ta}_4$ , shift to higher BE, while atoms of higher coordination, for example, central atoms in a  $\text{Ta}_{13}$  cluster, show more bulk-like behavior and shift to lower BE. This results in photoemission broadening at lower BE with increasing cluster size. The Bader charge displays a similar trend with BE and correlates with the Ta–Pt distance. Ta atoms far from the Pt surface display a low charge compared to a high charge for Ta atoms close to the surface. Both an increase in CN and Ta–Pt distance reduces the Bader charge of the Ta atom in the cluster, thereby influencing its oxidation state. At 40 K, all clusters are stable, and no beam damage is observed. In comparison, at 300 K,  $\text{Ta}_8$  is stable, as confirmed by STM, but  $\text{Ta}_4$  clusters show fragmentation induced by a thermodynamic drive for dispersion that the DFT calculations can rationalize. Notably, the stability of clusters larger than  $\text{Ta}_8$  can arguably be expected under these conditions. Furthermore, simulation and STM indicate an overall flat geometry of  $\text{Ta}_4$  and  $\text{Ta}_8$  clusters adsorbed to the Pt(111) surface, and based on the favorable

interaction with Pt, this flat structure will hold for all  $\text{Ta} < 8$  clusters at least. The origin of emissions concomitant to Ta adatoms in the cluster XPS spectra at 40 K is elucidated and reasoned to be related to the oxidation of Ta clusters. Slight traces of oxygen are sufficient to induce considerable changes to the XPS spectra that indicate the formation of Ta-oxides. The oxidation mechanism of the  $\text{Ta}_8$  clusters was treated computationally and appears to result via the formation of Ta-oxo-clusters with a concomitant abstraction of Ta atoms. A subsequent agglomeration of arguably mobile Ta-oxo-clusters of reduced size and Ta adatoms yields the final product of oxidation. STM combined with simulation reveals predominantly flat Ta-oxide islands, regardless of the cluster size, that form upon oxidation. At the low temperatures where it occurs, the process is not thermally but oxidatively driven. Spectroscopically, the Ta-oxide islands display a peculiar CLS inconsistent with values found in the literature for fully oxidized  $\text{Ta}^{5+}$ . Complementary DFT calculations indicate distinct surface and interfacial Ta binding configurations in the Ta-oxide islands on Pt(111). Their BEs are strongly affected by the Pt-induced perturbation, which effectively reduces the observed BEs with respect to bulk  $\text{Ta}_2\text{O}_5$ , most prominently for Ta atoms in the vicinity of Pt. This outcome underlines the importance of fundamental insights into the interfacial interaction between Ta and Pt. The experimentally observed CLSs, together with the computational findings, emphasize that the measured Ta 4f BEs cannot be solely explained in terms of the Ta oxidation state. Last but not least, we demonstrate the environmentally driven duality between surface segregation and subsurface intermixing of Ta on the Pt support. We observe a strong tendency of Ta to segregate in the presence of oxygen to form Ta-oxide islands, while reductive annealing in UHV (>900 K) leads to intermixing of Ta into the Pt subsurface. Cyclic interconversion between the two states can be achieved and offers promising bimetallic material properties for potential application.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jpcc.5c00699>.

Simulated gas-phase cluster geometry (Figure S1),  $\text{Ta}_4/\text{Pt}(111)$  STM image (Figure S2), first vs last Ta 4f XP scan all sizes (Figure S3), Ta 4f XP scans oxidation all sizes (Figure S4), O 1s/O 2s XP scans Pt(111) blanks (Figure S5), simulated oxidation mechanisms (Figure S6), Ta 4f XP scans (10 L) all sizes (Figure S7), O 1s/O 2s XP scan (10 L) references (Figure S8), Pt 4f XP scans Pt(111) blanks (Figure S9), computational values  $\text{Ta}_4$  (Table S1), computational values  $\text{Ta}_8$  (Table S2), Ta–Ta distances  $\text{str}_1(\text{Ta}_8)/\text{Pt}(111)$  (Table S3), Ta–Ta distances  $\text{Ta}_8$  gas-phase (Table S4), computational values Ta oxidation (Table S5), Ta–Ta distances  $\text{str}_2(\text{Ta}_8)/\text{Pt}(111)$  (Table S6) (PDF)

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## Notes

The authors declare no competing financial interest.

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## ABBREVIATIONS

XPS, X-ray photoelectron spectroscopy; HR-XPS, high-resolution X-ray photoelectron spectroscopy; DFT, density functional theory; STM, scanning tunneling microscopy; UHV, ultrahigh vacuum; CLS, core-level shift; IER, internal experimental reference; BE, binding energy;  $E_{\text{ads}}$ , adsorption energy; str, structure; CN, in-plane coordination number;  $\text{O}_{\text{chem}}$ , chemisorbed molecular oxygen

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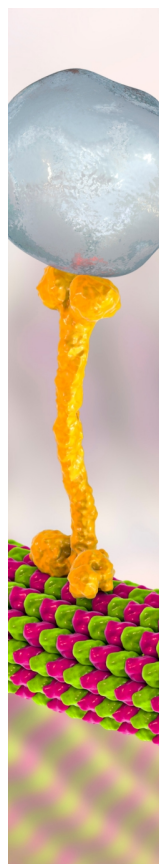
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